

湮灭算符:  $\hat{a} = (\frac{m\omega}{2\hbar})^{\frac{1}{2}}(\hat{x} + \frac{i}{m\omega}\hat{p})$

产生算符:  $\hat{a}^\dagger = (\frac{m\omega}{2\hbar})^{\frac{1}{2}}(\hat{x} - \frac{i}{m\omega}\hat{p})$

证明对易关系:  $[\hat{a}, \hat{a}^\dagger] = 1$

$$\begin{aligned} [\hat{a}, \hat{a}^\dagger] &= [(\frac{m\omega}{2\hbar})^{\frac{1}{2}}(\hat{x} + \frac{i}{m\omega}\hat{p}), (\frac{m\omega}{2\hbar})^{\frac{1}{2}}(\hat{x} - \frac{i}{m\omega}\hat{p})] \\ &= \frac{m\hbar}{2\hbar}[\hat{x} + \frac{i}{m\omega}\hat{p}, \hat{x} - \frac{i}{m\omega}\hat{p}] \\ &= \frac{m\hbar}{2\hbar}\{[\hat{x}, \hat{x}] - [\hat{x}, \frac{i}{m\omega}\hat{p}] + [\frac{i}{m\omega}\hat{p}, \hat{x}] - [\frac{i}{m\omega}\hat{p}, \frac{i}{m\omega}\hat{p}]\} \\ &= \frac{m\hbar}{2\hbar}\{-\frac{i}{m\omega}[\hat{x}, \hat{p}] + \frac{i}{m\omega}[\hat{p}, \hat{x}]\} \\ &= \frac{m\hbar}{2\hbar} \cdot \frac{2\hbar}{m\omega} \\ &= 1 \end{aligned}$$

对于一维线性谐振子:

$$\xi = \sqrt{\frac{m\omega}{\hbar}}x = \alpha x$$

从而得出:

$$\hat{a} = \frac{1}{\sqrt{2}}(\xi + \frac{\partial}{\partial\xi})$$

$$\hat{a}^\dagger = \frac{1}{\sqrt{2}}(\xi - \frac{\partial}{\partial\xi})$$

将 $\hat{a}$ 作用于谐振子哈密顿算符的第 $n$ 个本征态 $\psi_n$ , 可以得到:

$$\hat{a}\psi_n = \frac{1}{\sqrt{2}}(\xi + \frac{\partial}{\partial\xi})\psi_n$$

再根据:

$$\xi\psi_n = \sqrt{\frac{n}{2}}\psi_{n-1} + \sqrt{\frac{n+1}{2}}\psi_{n+1}$$

$$\frac{d}{d\xi}\psi_n = \sqrt{\frac{n}{2}}\psi_{n-1} - \sqrt{\frac{n+1}{2}}\psi_{n+1}$$

最终得到:

$$\begin{aligned}
\hat{a}\psi_n &= \frac{1}{\sqrt{2}}\xi\psi_n + \frac{1}{\sqrt{2}}\frac{\partial}{\partial\xi}\psi_n \\
&= \frac{1}{\sqrt{2}}\left(\sqrt{\frac{n}{2}}\psi_{n-1} + \sqrt{\frac{n+1}{2}}\psi_{n+1} + \sqrt{\frac{n}{2}}\psi_{n-1} - \sqrt{\frac{n+1}{2}}\psi_{n+1}\right) \\
&= \frac{1}{\sqrt{2}} \cdot 2 \cdot \sqrt{\frac{n}{2}}\psi_{n-1} \\
&= \sqrt{n}\psi_{n-1}
\end{aligned}$$

同理：

$$\hat{a}^\dagger\psi_n = \sqrt{n+1}\psi_{n+1}$$